

## The right result for the ~~wrong~~ RIGHT reason!

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Radisson SAS Royal Hotel Copenhagen, April 29, 2006

### Software

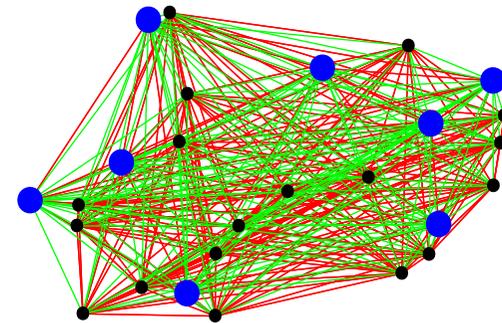
- Dalton (<http://www.kjemi.uio.no/software/dalton>)
  - LUCIA (J. Olsen)
  - ACES II (J. Gauss, J. Stanton, R. Bartlett)

## High-precision quantum-chemical calculations

- The last decade has seen a dramatic improvement in our ability to treat molecular electronic systems accurately
  - development of techniques for systematic convergence towards the exact solution
  - extensive benchmarking on small and light molecular systems
- In many cases, we can now confidently confirm or reject experimental observations
- Many black-box methods have been developed
  - well-defined levels of theory
- Still, the exact solution can be approached in infinitely many ways
  - at many (incomplete) levels of theory, agreement with experiment may be obtained
  - such agreement always arises by error cancellation
  - error cancellation is treacherous: the right answer for the wrong reason
- For a reliable application of quantum-chemistry methods, one must be wary of the pitfalls of error cancellation
- We shall here review high-precision quantum chemistry, paying special attention to the problem of error cancellation

## The Hartree–Fock model

- All information about the electronic system is contained in the **wave function**, obtained by solving the **Schrödinger equation** (or the corresponding relativistic equation).
  - Molecular electronic systems represent a difficult **many-body problem** that cannot be treated exactly—even on modern computers.
  - We are thus forced to make **approximations**, setting up simplified descriptions that incorporate the most important features of the system.
- This should preferably be done in an orderly fashion, so that the exact solution can be approached in a systematic manner: **hierarchy of approximations**
- **The Hartree–Fock model**—the fundamental approximation of wave-function theory
  - each electron moves in the **mean field** of all other electrons
  - provides an **uncorrelated description**: average rather than instantaneous interactions
  - gives rise to the concept of **molecular orbitals**
  - **typical errors**: 0.5% in the energy; 1% in bond distances, 5%–10% in other properties
  - forms the basis for more accurate treatments



## Electron correlation and virtual excitations

- **electron correlation:**
  - to improve upon the HF model, we must take into account the instantaneous interactions among the electrons
  - in real space, the electrons are constantly being scattered by **collisions**
  - in the orbital picture, these collisions manifest themselves as **excitations** from occupied to virtual (unoccupied) spin orbitals
- **double excitations:**
  - the most important events are collisions between two electrons
  - in the orbital picture, such an event corresponds to an excitation from two occupied to two virtual spin orbitals, known as **pair excitations** or **double excitations**
  - with each such excitation, we associate an **amplitude**, representing its probability
  - our description is then said to be **correlated**
- **coupled-cluster theory:**
  - by allowing all possible single and double excitations to happen in all possible combinations, we arrive at the **coupled-cluster singles-and-doubles (CCSD)** model
  - in the CCSD model, the HF error is typically reduced by a factor of three or four
  - for higher accuracy, we must also include triple excitations (CCSDT), and so on

## Basis sets of Gaussian functions

- In our calculations, we expand the molecular orbitals in one-electron Gaussian-type functions (GTOs):

$$G_{ijk}(r_A, \alpha) = x_A^i y_A^j z_A^k \exp(-\alpha r_A^2)$$

- Basis sets of increasing size:
  - minimal or single-zeta (SZ) basis sets:
    - \* one set of GTOs for each occupied atomic shell ( $2s1p$ )
    - \* gives a rudimentary description of electron structure
  - double-zeta (DZ) basis sets:
    - \* two sets of GTOs for each occupied atomic shell ( $3s2p1d$ )
    - \* sufficient for a qualitative description of the electron system
  - triple-zeta (TZ), quadruple-zeta (QZ) and larger basis sets:
    - \* needed for a quantitative description of the electronic system
- The number of GTOs per atom increases rapidly:

SZ	DZ	TZ	QZ	5Z	6Z
5	14	30	55	91	140

- Large basis sets are needed to generate a flexible virtual space!

## The two-dimensional chart of nonrelativistic quantum chemistry

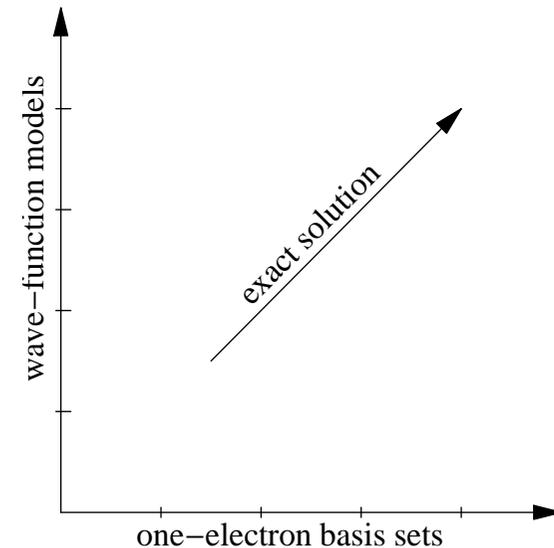
- The quality of nonrelativistic molecular electronic-structure calculations is determined by the description of
  1. the  $N$ -electron space (wave-function model),
  2. the one-electron space (basis set).
- In each space, there is a hierarchy of levels of increasing complexity:

1. the  $N$ -electron hierarchy:  
coupled-cluster excitation levels

HF, CCSD, CCSDT, CCSDTQ, ...

2. the one-electron hierarchy:  
correlation-consistent basis sets

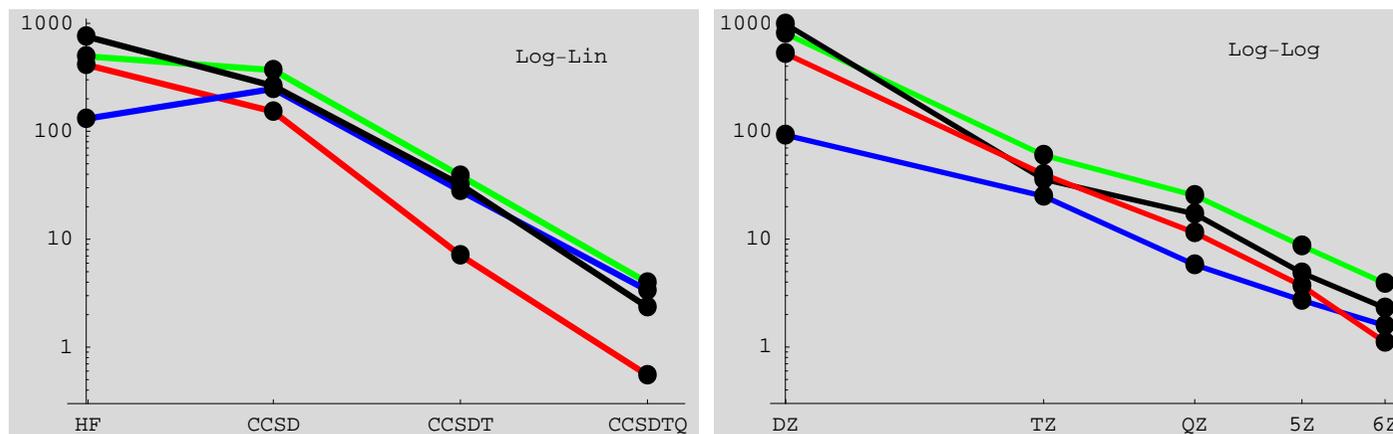
DZ, TZ, QZ, 5Z, 6Z, ...



- The quality is systematically improved upon by going up in the hierarchies.

## Energy contributions to atomization energies (kJ/mol)

- Contributions of each CC excitation level (left) and AO basis-set shell (right)



– color code: **HF**, **N<sub>2</sub>**, **F<sub>2</sub>**, and **CO**

- The excitation-level convergence is approximately linear (log-linear plot)
  - each new excitation level reduces the error by about an order of magnitude
  - the contributions from quintuples are negligible (about 0.1 kJ/mol)
- The basis-set convergence is much slower (log-log plot)
  - each shell contributes an energy proportional to  $X^{-4}$  where  $X$  is the cardinal number
  - a similarly small error (0.1 kJ/mol) requires  $X > 10$
  - clearly, we must choose our orbitals in the best possible manner

## Contributions to atomization energies (kJ/mol)

	RHF	SD	T	Q	rel.	vib.	total	experiment	error
CH <sub>2</sub>	531.1	218.3	9.5	0.4	-0.7	-43.2	715.4	714.8±1.8	0.6
H <sub>2</sub> O	652.3	305.3	17.3	0.8	-2.1	-55.4	918.2	917.8±0.2	0.4
HF	405.7	178.2	9.1	0.6	-2.5	-24.5	566.7	566.2±0.7	0.5
N <sub>2</sub>	482.9	426.0	42.4	3.9	-0.6	-14.1	940.6	941.6±0.2	-1.1
F <sub>2</sub>	-155.3	283.3	31.6	3.3	-3.3	-5.5	154.1	154.6±0.6	-0.5
CO	730.1	322.2	32.1	2.3	-2.0	-12.9	1071.8	1071.8±0.5	-0.0

- agreement with experiment in all cases except for N<sub>2</sub>
- all purely electronic contributions are positive (except Hartree–Fock for F<sub>2</sub>)
  - very large correlation contributions
  - the neglected quintuples contribute a few tenths of 1 kJ/mol
- relativistic corrections are small ( $\approx 0.5\%$ ) but needed for agreement with experiment
  - nearly cancel quadruples contributions
- vibrational corrections are substantial
  - similar to triples contributions but oppositely directed

## Contributions to bond distances (pm)

	RHF	SD	T	Q	5	rel.	theory	exp.	err.
HF	89.70	1.67	0.29	0.02	0.00	0.01	91.69	91.69	0.00
N <sub>2</sub>	106.54	2.40	0.67	0.14	0.03	0.00	109.78	109.77	0.01
F <sub>2</sub>	132.64	6.04	2.02	0.44	0.03	0.05	141.22	141.27	-0.05
CO	110.18	1.87	0.75	0.04	0.00	0.00	112.84	112.84	0.00

- agreement with experiment to within 0.01 pm except for F<sub>2</sub>
- Hartree–Fock theory underestimates bond distances by up to 8.6 pm (for F<sub>2</sub>)
- all correlation contributions are positive
  - approximate linear convergence, slowest for F<sub>2</sub>
  - triples contribute up to 2.0 pm, quadruples up to 0.4 pm, and quintuples 0.03 pm
  - sextuples are needed for convergence to within 0.01 pm
- relativistic corrections are small except for F<sub>2</sub> (0.05 pm)
  - of the same magnitude and direction as the quintuples

## Contributions to harmonic frequencies $\omega_e$ ( $\text{cm}^{-1}$ )

	RHF	SD	T	Q	5	rel.	theory	exp.	err.
HF	4473.8	-277.4	-50.2	-4.1	-0.1	-3.5	4138.5	4138.3	0.2
N <sub>2</sub>	2730.3	-275.8	-72.4	-18.8	-3.9	-1.4	2358.0	2358.6	-0.6
F <sub>2</sub>	1266.9	-236.1	-95.3	-15.3	-0.8	-0.5	918.9	916.6	2.3
CO	2426.7	-177.4	-71.7	-7.2	0.0	-1.3	2169.1	2169.8	0.7

- agreement with experiment to within  $1 \text{ cm}^{-1}$  except for F<sub>2</sub>
- Hartree–Fock theory overestimates harmonic frequencies by up to 38% (in F<sub>2</sub>).
- all correlation contributions are large and negative
  - triples contribute up to  $95 \text{ cm}^{-1}$ , quadruples  $20 \text{ cm}^{-1}$ , and quintuples  $4 \text{ cm}^{-1}$
  - sextuples are sometimes needed for convergence to within  $1 \text{ cm}^{-1}$
- relativistic corrections are of the order of  $1 \text{ cm}^{-1}$ 
  - of the same magnitude and direction as the quadruples or quintuples

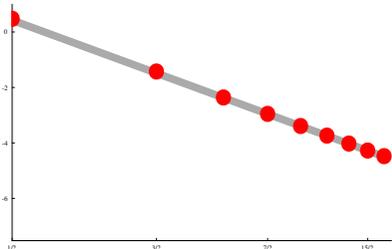
# The principal expansion and correlation-consistent basis sets

- The energy contribution from each AO in large CI calculations on helium:

$$\varepsilon_{nlm} \approx n^{-6} \quad \leftarrow \text{Carroll et al. (1979)}$$

- The principal expansion:** include all AOs belonging to the same shell simultaneously, in order of increasing principal quantum number  $n$ :

$$\varepsilon_n \approx n^2 n^{-6} = n^{-4}$$



- Practical realization: the correlation-consistent basis sets **cc-pVXZ (Dunning, 1989)**
- Energy-optimized AOs are added one shell at a time:

SZ	cc-pVDZ	cc-pVTZ	cc-pVQZ	number of AOs
	+3s3p3d	+4s4p4d4f	+5s5p5d5f5g	$\propto X^2$
2s1p	3s2p1d	4s3p2d1f	5s4p3d2f1g	$\propto X^3$

- The error in the energy is equal to the contributions from all omitted shells:

$$\Delta E_X \approx \sum_{n=X+1}^{\infty} n^{-4} \approx X^{-3} \approx N^{-1} \approx T^{-1/4}$$

- Each new digit in the energy therefore costs 10000 times more CPU time!

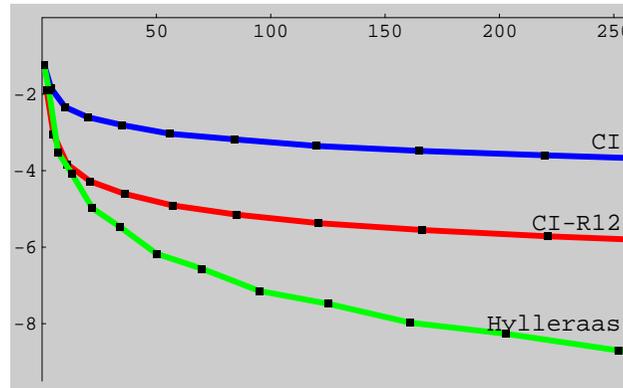
$$1 \text{ minute} \quad \rightarrow \quad 1 \text{ week} \quad \rightarrow \quad 200 \text{ years}$$

## Solutions to slow basis-set convergence

### 1. Use explicitly correlated methods!

- Include interelectronic distances  $r_{ij}$  in the wave function (Hylleraas 1928):

$$\Psi_{R12} = \sum_K C_K \Phi_K + C_{R12} r_{12} \Phi_0$$



- We use CCSD-R12 (Noga, Kutzelnigg, and Klopper, 1992) for benchmarking
- Note: basis-set convergence mainly a problem for double excitations

### 2. Use basis-set extrapolation!

- Exploit the smooth convergence  $E_\infty = E_X + AX^{-3}$  to extrapolate to basis-set limit

$$E_\infty = \frac{X^3 E_X - Y^3 E_Y}{X^3 - Y^3}$$

$mE_h$	DZ	TZ	QZ	5Z	6Z	R12
plain	194.8	62.2	23.1	10.6	6.6	1.4
extr.		21.4	1.4	0.4	0.5	

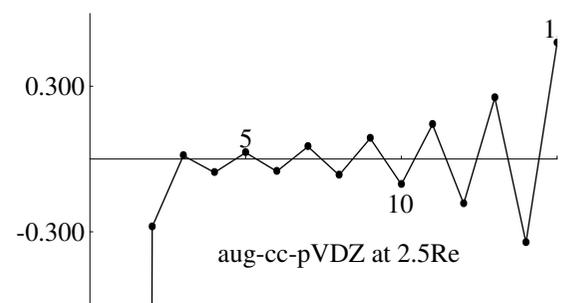
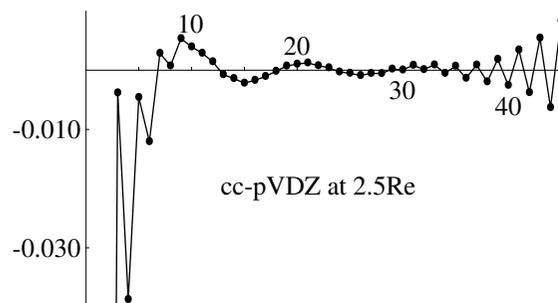
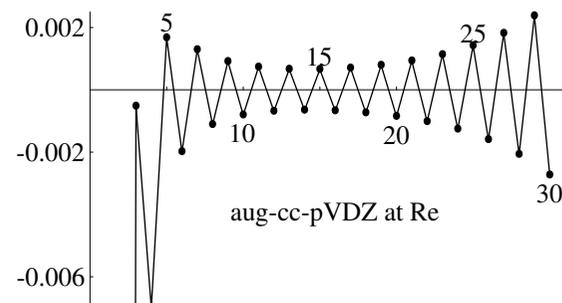
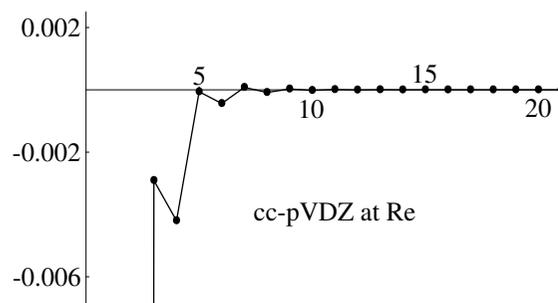
- The formula is linear and contains no parameters; applicable to many properties

## Practical levels of theory: error cancellation

- We have seen how the “exact” result can be approached systematically.
- In practice, we cannot carry out such studies in chemical applications
  - only relatively low excitation levels can be reached
  - only relatively small basis sets can be used
- To reduce cost, additional (simplified) levels are introduced
  - typically based on perturbation theory: MP2, CCSD(T)
  - introduces additional sources of error cancellations
- We shall now examine the convergence more carefully
  - emphasis on the interplay between excitation-level and basis-set convergence
- We shall do this for the three properties considered about:
  - atomization energies
  - bond distances
  - vibrational frequencies

## Many-body perturbation theory: approximate coupled-cluster theory

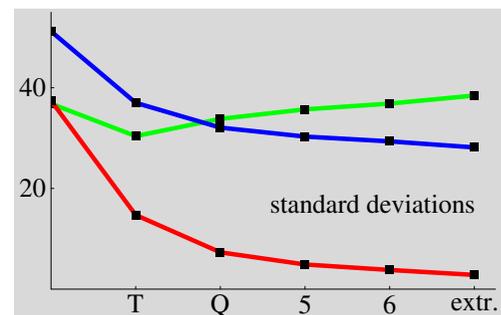
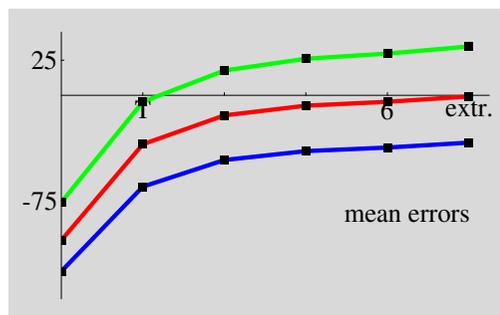
- Coupled-cluster amplitudes may be estimated by **perturbation theory**
- Caveat: the resulting perturbation series is frequently divergent, even in simple cases
- Here are some examples for the HF molecule (10 electrons):



- However, to lowest order, perturbational corrections are very useful and popular
  - **MP2** (approximate CCSD) and **CCSD(T)** (approximate CCSDT)

## Atomization energies (AEs)

- Let us consider the situation for an important molecular property: AEs
- Statistics based 20 closed-shell organic molecules (kJ/mol)



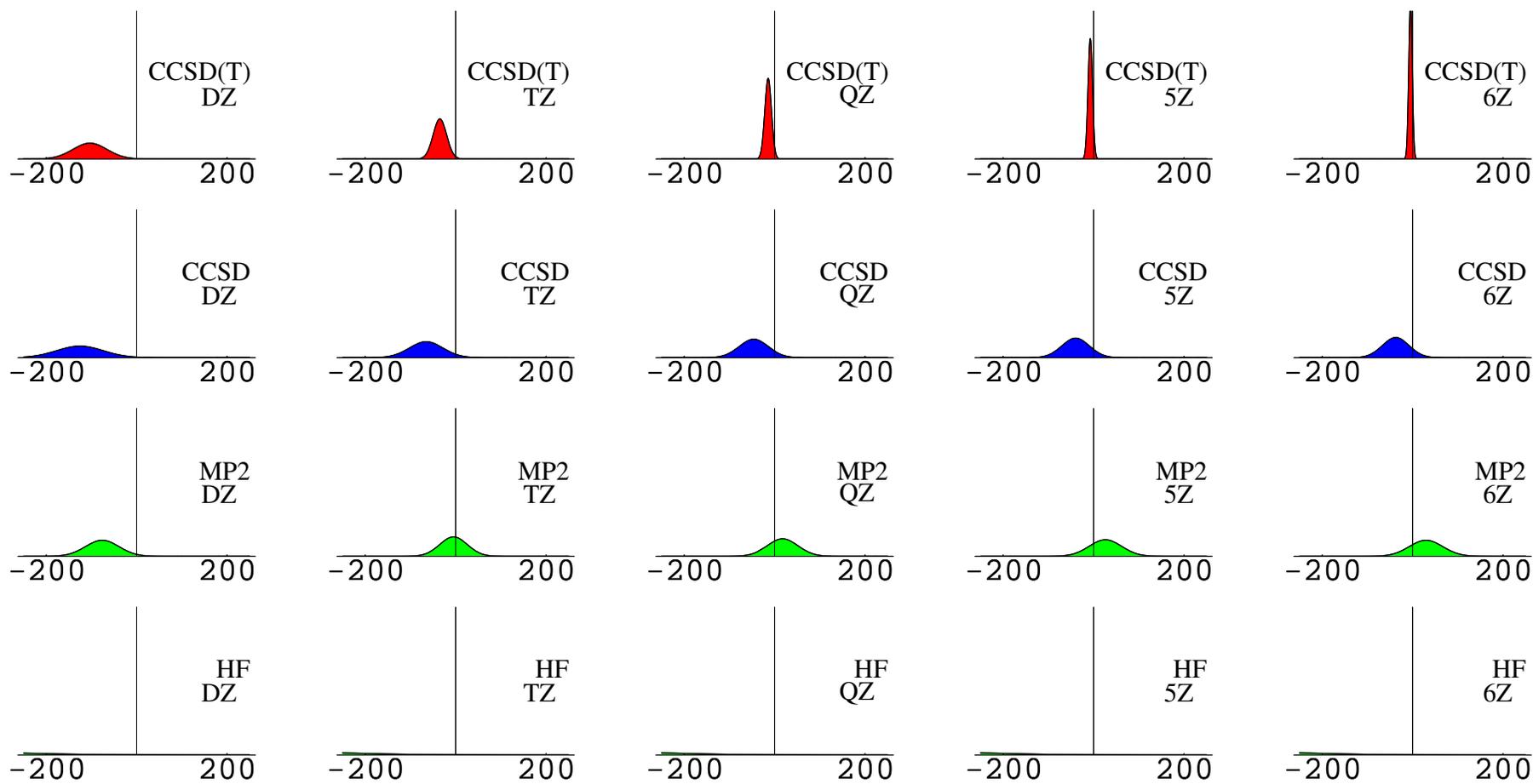
- AEs increase with excitation level in the coupled-cluster hierarchy:

$$\text{HF} < \text{CCSD} < \text{CCSD(T)} < \text{MP2}$$

- Mean abs. cc-pcV6Z err. (kJ/mol): 423 (HF), 37 (MP2), 30 (CCSD), 4 (CCSD(T))
- AEs increase with cardinal number.
- CCSD(T) performs excellently, but DZ and TZ are inadequate:

kJ/mol	DZ	TZ	QZ	5Z	6Z
raw	103.1	34.0	13.5	6.6	4.1
extrapolated		14.5	1.7	0.9	0.8

# Atomization energies (kJ/mol)



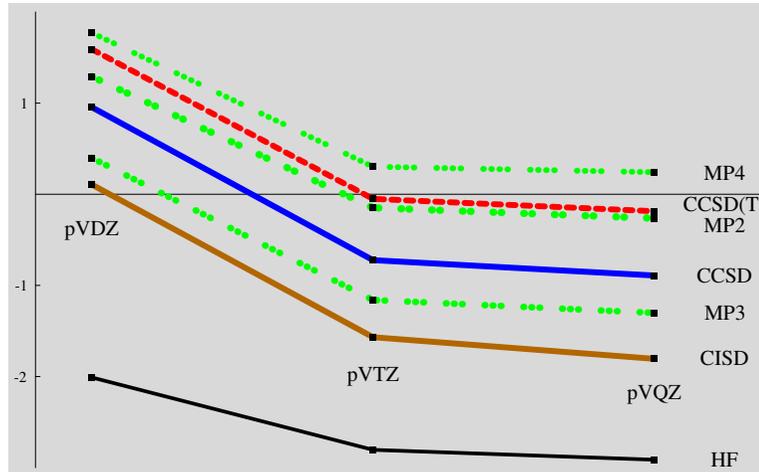
## The (in)adequacy of CCSD(T)

	CCSD(T)		CCSDT		CCSDTQ		experiment	
	cc-pCV(56)Z		cc-pCV(Q5)Z		cc-pVTZ		$D_e$	$D_0$
CH <sub>2</sub>	757.9	-0.9	758.9	0.1	759.3	0.5	758.8	714.8±1.8
H <sub>2</sub> O	975.3	0.1	974.9	-0.3	975.7	0.5	975.2	917.8±0.2
HF	593.2	0.0	593.0	-0.2	593.6	0.4	593.2	566.2±0.7
N <sub>2</sub>	954.7	-1.6	951.3	-5.0	955.2	-1.1	956.3	941.6±0.2
F <sub>2</sub>	161.0	-2.4	159.6	-3.8	162.9	-0.5	163.4	154.6±0.6
CO	1086.7	0.0	1084.4	-2.3	1086.7	0.0	1086.7	1071.8±0.5

- The excellent performance of CCSD(T) for AEs relies on error cancellation:
  - relaxation of triples from CCSD(T) to CCSDT **reduces** the AEs;
  - inclusion of quadruples from CCSDT to CCSDTQ **increases** the AEs.
- The error incurred by treating the connected triples perturbatively is quite large (about 10% of the full triples contribution) but canceled by the neglect of quadruples.
- The rigorous calculation of AEs to chemical accuracy requires **CCSDTQ/cc-pCV6Z!**

## Bond distances: statistics

- Statistics based on 28 bond distances at the all-electron cc-pVXZ level (pm):



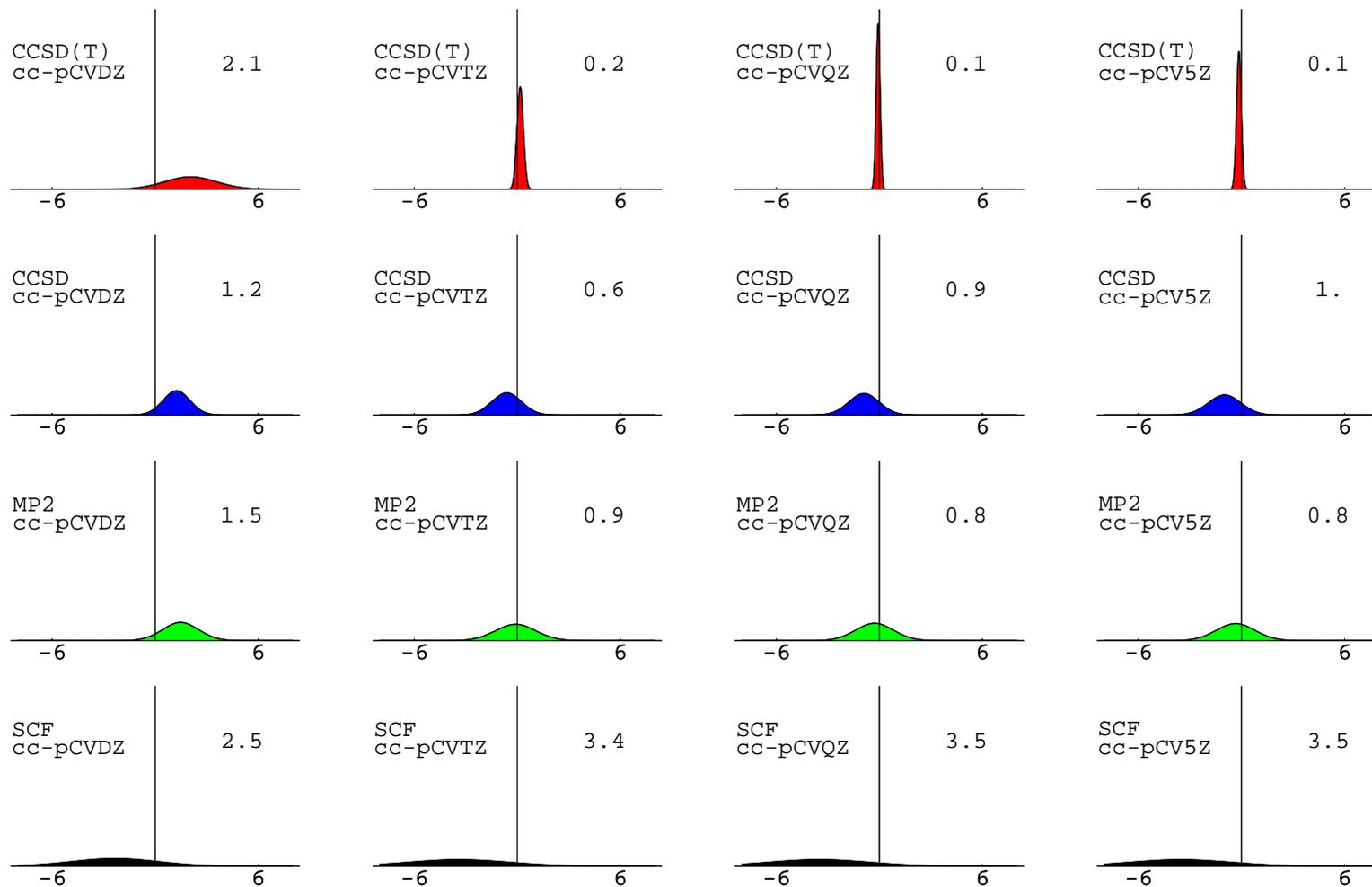
- Bonds shorten with increasing basis:  
DZ > TZ > QZ
- Bonds lengthen with increasing excitations:  
HF < CCSD < MP2 < CCSD(T)
- Considerable scope for error cancellation:  
CISD/DZ, MP3/DZ

- CCSD(T) distances compared with experiment:

pm	DZ	TZ	QZ
$\bar{\Delta}$	1.68	0.01	-0.12
$ \bar{\Delta} $	1.68	0.20	0.16

- The high accuracy of CCSD(T) arises partly because of error cancellation.
- Bond distances are further reduced by
  - basis-set extension QZ  $\rightarrow$  6Z:  $\approx -0.10$  pm
  - triples relaxation CCSD(T)  $\rightarrow$  CCSDT:  $\approx -0.04$  pm
- Intrinsic error of the CCSDT model:  $\approx -0.2$  pm
- Connected quadruples increase bond lengths by about 0.1–0.2 pm

# Bond distances $R_e$ of BH, CO, N<sub>2</sub>, HF, and F<sub>2</sub> (cm<sup>-1</sup>)



## Error cancellation: some general considerations

- In an infinite basis, the CC hierarchy is related to experiment as follows:
  - $\text{HF} < \text{CCSD} < \text{CCSDT} < \text{CCSDTQ} < \dots < \text{FCI} = \text{exp}$  (BDs and AEs)
  - no error cancellation may occur (ignoring relativity and adiabatic corrections)
- In a finite basis, results are shifted and cancellation occurs for BDs but not for AEs:
  - $\text{HF} < \text{CCSD} \approx \text{exp} < \text{CCSDT} < \text{CCSDTQ} < \dots < \text{FCI}$  (BDs shifted up)
  - $\text{HF} < \text{CCSD} < \text{CCSDT} < \text{CCSDTQ} < \dots < \text{FCI} < \text{exp}$  (AEs shifted down)
- Perturbative treatments of connected excitations tend to overestimate their effect:
  - $\text{HF} < \text{CCSD} < \text{MP2} \leftarrow$  (over)compensates for lack of triples
  - $\text{CCSD} < \text{CCSDT} < \text{CCSD(T)} \leftarrow$  compensates for lack of quadruples
  - a new source of error cancellation
- For BDs and AEs, we now observe
  - $\text{HF} < \text{CCSD} < \text{MP2} \approx \text{CCSD(T)} \approx \text{exp}$  (BDs)
  - $\text{HF} < \text{CCSD} < \text{CCSD(T)} \approx \text{exp} < \text{MP2}$  (AEs)
  - error cancellation may now occur also for AEs
- We shall see that vibrational frequencies behave in much the same manner as BDs since basis-set extension and increased excitation levels work in opposite directions

## Vibrational frequencies of diatoms

- The frequency of a diatomic fundamental transition is given by

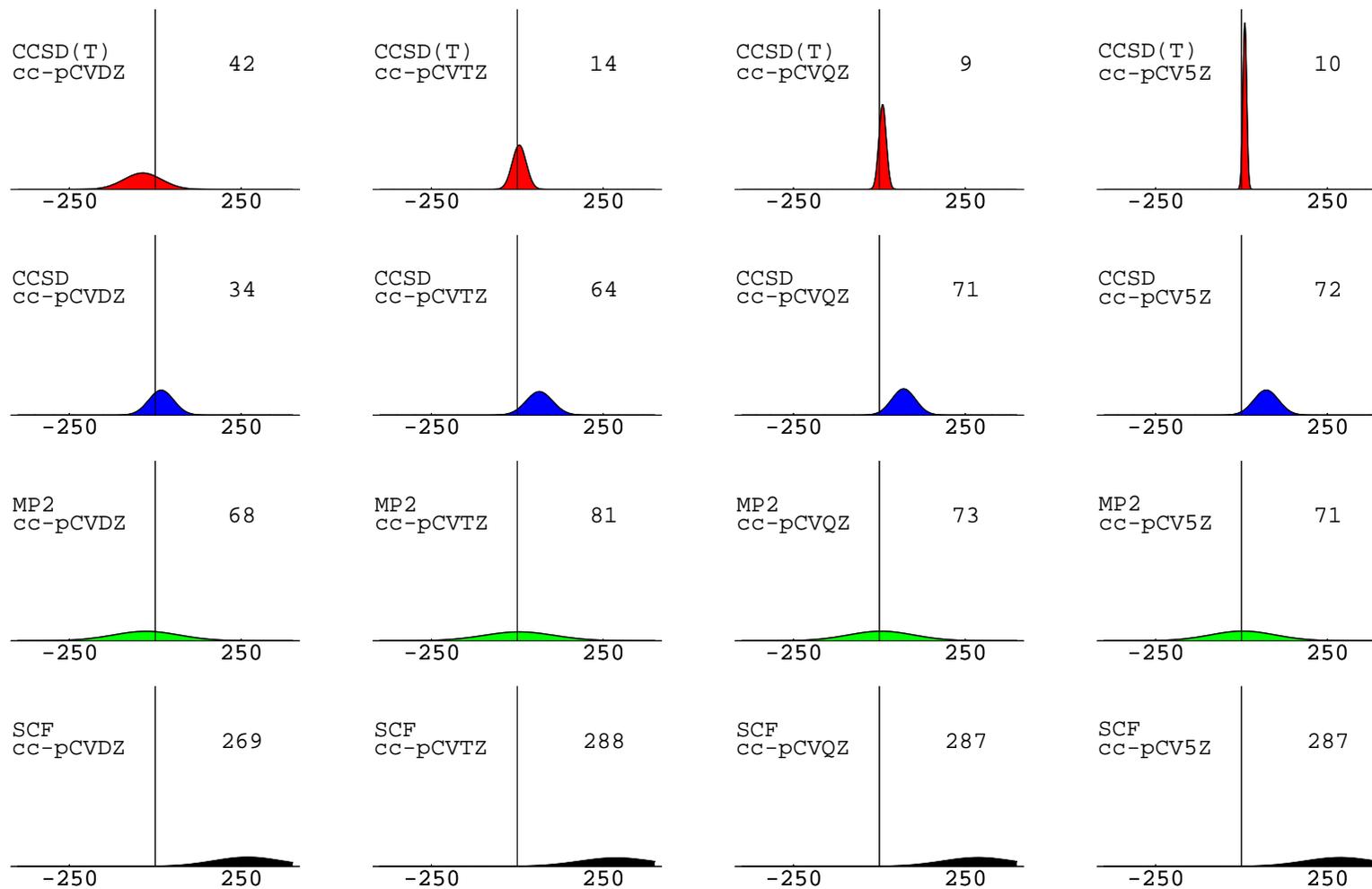
$$\nu = \omega_e - 2\omega_e x_e$$

- the harmonic constant  $\omega_e$  requires 2nd derivatives of PES
- the anharmonic constant  $\omega_e x_e$  requires 4th derivatives of PES
- With the advent of CCSD(T) in the 1990s, it was soon realized that this model is capable of highly accurate vibrational constants (to within a few wavenumbers at the TZ level)
- However, with the development of codes capable of handling very large basis sets and high excitation levels, it has slowly transpired that things are perhaps not so simple
- It is easier to converge harmonic than anharmonic constants:

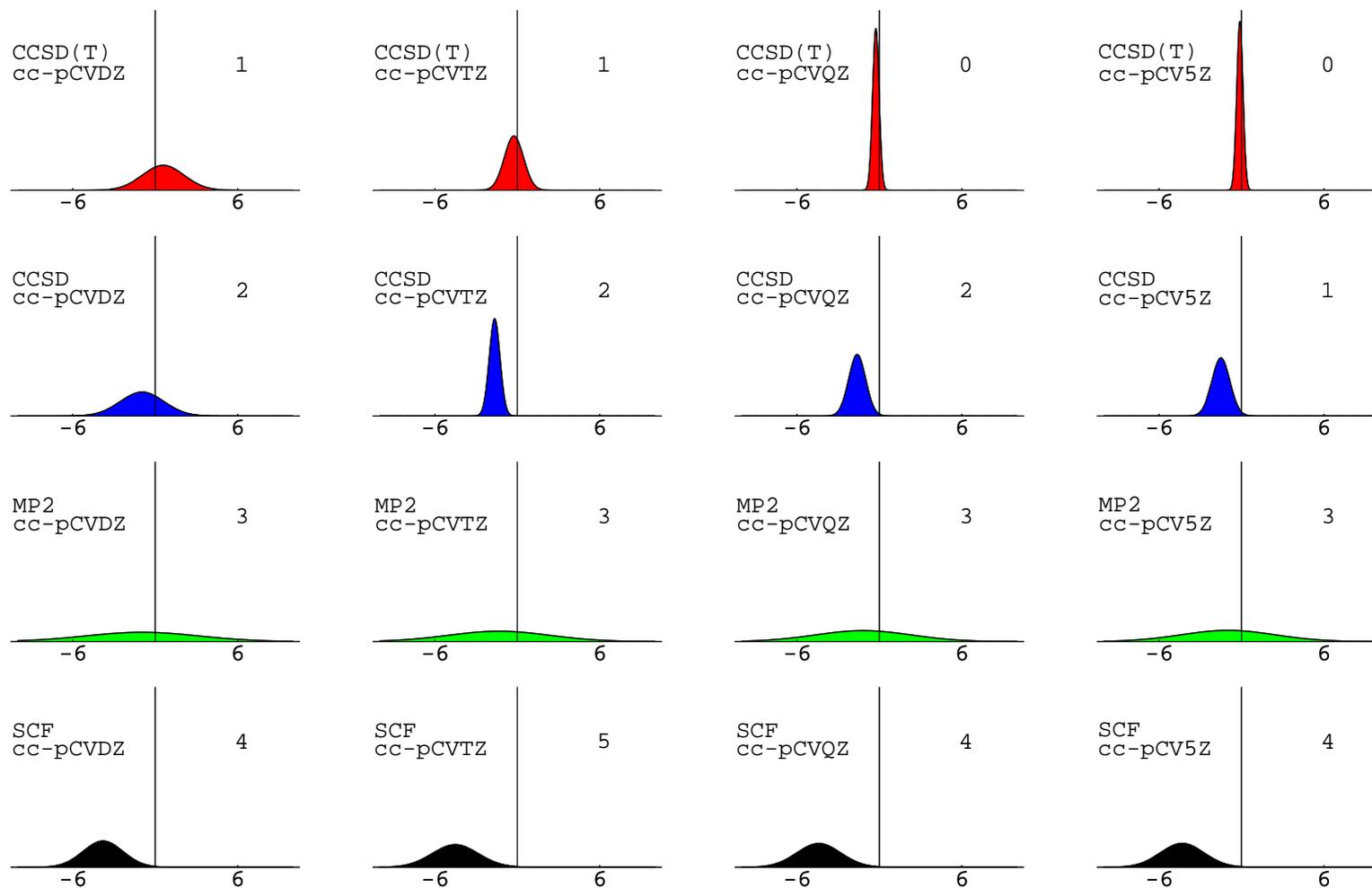
$ \overline{\Delta} $ (cm <sup>-1</sup> )	RHF	MP2	CCSD	CCSD(T)	judgment
$\omega_e$	287	71	72	10	DIFFICULT!
$\omega_e x_e$	4	3	1	0	EASY!

- There are two main problems with harmonic constants:
  - basis-set incompleteness
  - lack of high-order connected excitations

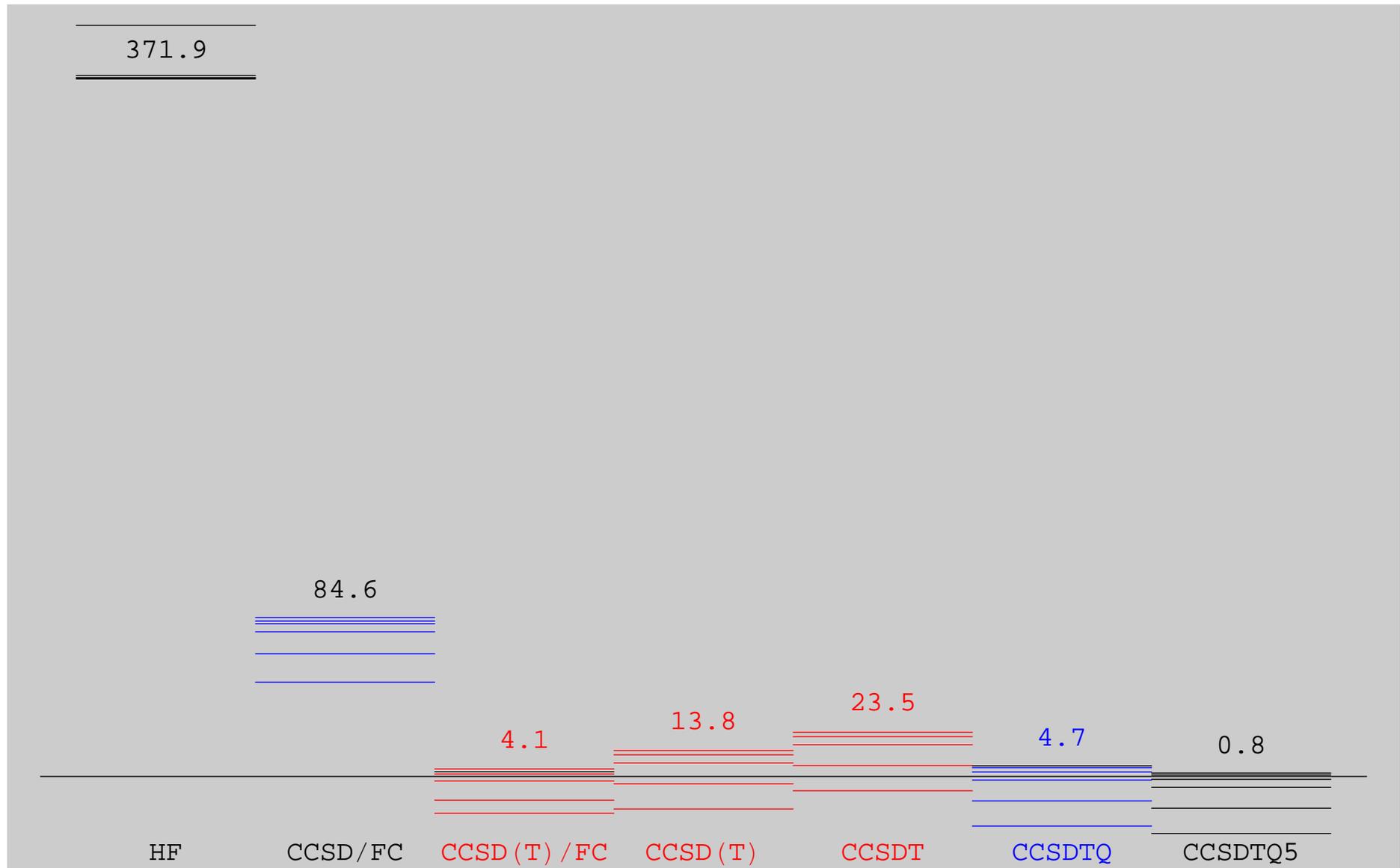
# Harmonic constants $\omega_e$ of BH, CO, N<sub>2</sub>, HF, and F<sub>2</sub> (cm<sup>-1</sup>)



# Anharmonic constants $\omega_e x_e$ of BH, CO, N<sub>2</sub>, HF, and F<sub>2</sub> (cm<sup>-1</sup>)



# Higher-order connected contributions to $\omega_e$ in $N_2$



## Conclusions

- Quantum chemistry provides a set of well-defined levels of approximation
  - excitation-level expansion
  - basis-set expansion
- The “exact” result can be approached in a systematic manner
- In practice, we have to be content with low levels of theory
  - low excitation levels, approximated by perturbation theory
  - small basis sets
- Sometimes complicated interplay between different approximations
  - useful and reliable (balanced) levels of theory have been implemented in codes
  - can be usefully applied without deep knowledge of quantum chemistry
- Still, some knowledge of these methods, their errors and the interplay of these errors, will make the application of these methods more reliable
- Discipline is always needed!